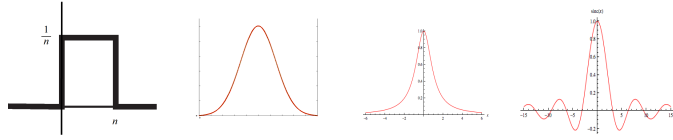


The Dirac Delta Function(al) $\delta(t)$



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The purpose of this document is to **illustrate the properties of Dirac Delta Function** $\int_{-\infty}^{\infty} \delta(x)dx = 1$

1 Definition

The Dirac Delta Function is *defined by its assigned properties*

1. It decays

$$\delta(x) = 0, x \neq 0$$

2. Screening property

$$\int_{-\infty}^{\infty} f(x)\delta(x)dx = f(0)$$

Where $f(x)$ is well-defined ordinary function

By property 2, if $f(x)$ is the unity function, i.e. $f(x) = 1(t) = 1$

$$\int_{-\infty}^{\infty} 1(x)\delta(x)dx = 1(0) = 1$$

Thus

$$\int_{-\infty}^{\infty} \delta(x)dx = 1$$

Therefore, the Dirac Delta function has the following properties

It is infinitely thin spike
It is infinitely high
It is not ordinary function

2 Dirac Delta Function as a Functional

Since there is no ordinary function has such properties, the Dirac Delta function has to be represented as a *limit* of certain function

That means, Dirac Delta Function has *several expressions / several forms*. And all these forms must fill the assigned properties (1) and (2)

In this way, the Dirac Delta Function is actually a *function of functions* called *generalized function* or *functional* or *distribution*

It is a *function of functions* because it describes the general behaviours of a set of functions.

Following are some candidates of Dirac Delta Function

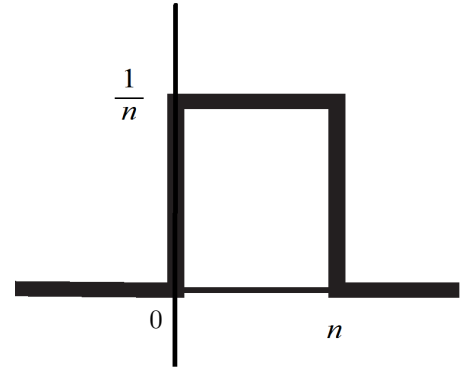
$$\delta_n(t) = \begin{cases} \frac{1}{n} & 0 \leq x \leq n \\ 0 & \text{else} \end{cases} \quad \delta_n(t) = \begin{cases} n & |x| \leq \frac{1}{2n} \\ 0 & \text{else} \end{cases} \quad \delta_n(t) = \sqrt{\frac{n^2}{\pi}} \exp(-n^2 x^2) \quad \delta_n(t) = \frac{n}{\pi} \frac{1}{1+n^2 x^2} \quad \delta_n(t) = \frac{\sin nx}{\pi x}$$

With these ordinary function representation, the Dirac Delta Function can be represented as

$$\int_{-\infty}^{\infty} f(x)\delta(x)dx = \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} f(x)\delta_n(x)dx = f(0)$$

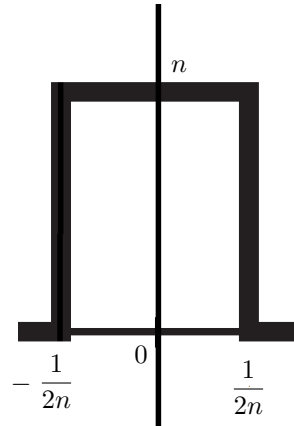
2.1 Rectangular Pulse

$$\begin{aligned} \delta_n(t) &= \begin{cases} \frac{1}{n} & 0 \leq x \leq n \\ 0 & \text{else} \end{cases} \\ &= \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} \delta_n(x) dx \\ &= \lim_{n \rightarrow \infty} \left[\underbrace{\int_{-\infty}^0 \delta_n(x) dx}_0 + \int_0^n \delta_n(x) dx + \underbrace{\int_n^{\infty} \delta_n(x) dx}_0 \right] \\ &= \lim_{n \rightarrow \infty} \int_0^n \frac{1}{n} dx \\ &= \lim_{n \rightarrow \infty} 1 \\ &= 1 \end{aligned}$$



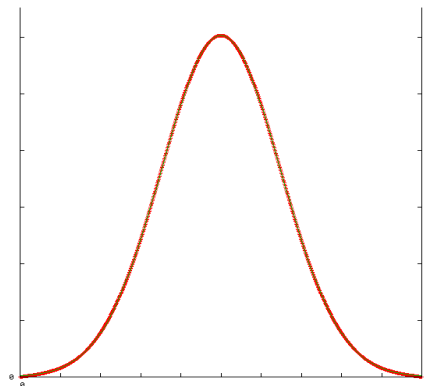
2.2 Central Rectangular pulse

$$\begin{aligned} \delta_n(t) &= \begin{cases} n & |x| \leq \frac{1}{2n} \\ 0 & \text{else} \end{cases} \\ &= \lim_{n \rightarrow \infty} \left(\int_{-\infty}^{-1/2n} \delta_n(x) dx + \int_{-1/2n}^{1/2n} \delta_n(x) dx + \int_{1/2n}^{\infty} \delta_n(x) dx \right) \\ &= \lim_{n \rightarrow \infty} n \int_{-1/2n}^{1/2n} dx \\ &= \lim_{n \rightarrow \infty} n \left[\frac{1}{2n} - \left(-\frac{1}{2n} \right) \right] \\ &= \lim_{n \rightarrow \infty} 1 \\ &= 1 \end{aligned}$$



2.3 Gaussian Pulse

$$\begin{aligned} \delta_n(t) &= \sqrt{\frac{n^2}{\pi}} \exp(-n^2 x^2) \\ &= \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} \sqrt{\frac{n^2}{\pi}} \exp(-n^2 x^2) dx \\ &= \lim_{n \rightarrow \infty} \sqrt{\frac{1}{\pi}} \int_{-\infty}^{\infty} \exp(-n^2 x^2) dx \\ &= \lim_{n \rightarrow \infty} \sqrt{\frac{1}{\pi}} \int_{-\infty}^{\infty} \exp(-y^2) dy \\ &= \lim_{n \rightarrow \infty} 2 \sqrt{\frac{1}{\pi}} \int_0^{\infty} \exp(-y^2) dy \end{aligned}$$



Using Jacobian Transformation technique

$$\int_0^{\infty} \exp(-y^2) dy = \frac{\sqrt{\pi}}{2}$$

The Jacobian Transformation Technique (For reference only !)

The following is an illustration of the Jacobian Transformation technique

$$\begin{aligned} & \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{(\int_0^{\infty} e^{-y^2} dy) (\int_0^{\infty} e^{-x^2} dx)}} \\ &= \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{(\int_0^{\infty} e^{-y^2} dy) (\int_0^{\infty} e^{-x^2} dx)}} \\ &= \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{\int_0^{\infty} \int_0^{\infty} e^{-(x^2+y^2)} dx dy}} \end{aligned}$$

From Rectangular $(x, y) \rightarrow$ Polar (r, θ)

$$\begin{cases} x = r \cos \theta \\ y = r \sin \theta \\ x^2 + y^2 = r^2 \\ \theta = \tan^{-1} \frac{y}{x} \end{cases} \quad \begin{cases} x \in [0, \infty) \\ y \in [0, \infty) \end{cases} \rightarrow \begin{cases} r \in [0, \infty) \\ \theta \in [0, \frac{\pi}{2}] \end{cases}$$

$$dx dy = J dr d\theta$$

Where the Jacobian

$$J = \frac{\partial(x, y)}{\partial(r, \theta)} = \det \begin{bmatrix} \frac{\partial x}{\partial r} & \frac{\partial x}{\partial \theta} \\ \frac{\partial y}{\partial r} & \frac{\partial y}{\partial \theta} \end{bmatrix} = \det \begin{bmatrix} \cos \theta & -r \sin \theta \\ \sin \theta & r \cos \theta \end{bmatrix} = r$$

Thus

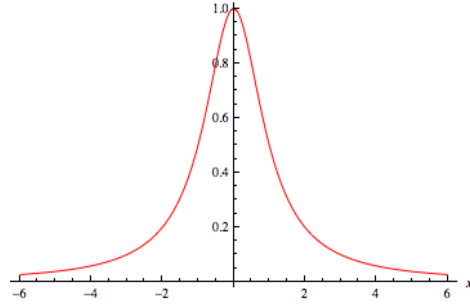
$$\begin{aligned} & \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{\int_0^{\infty} \int_0^{\infty} e^{-(x^2+y^2)} dx dy}} \\ &= \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{\int_0^{\pi/2} \int_0^{\infty} e^{-r^2} r dr d\theta}} \\ &= \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{\int_0^{\pi/2} d\theta \int_0^{\infty} e^{-r^2} r dr}} \\ &= \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{\frac{\pi}{2} \int_0^{\infty} e^{-r^2} \frac{1}{2} dr^2}} \\ &= \frac{\int_0^{\infty} e^{-y^2} dy}{\sqrt{\frac{\pi}{4} \int_0^{\infty} e^{-r^2} dr^2}} \\ &= \sqrt{\frac{\pi}{4}} \end{aligned}$$

Therefore

$$\begin{aligned} & \frac{\int_{-\infty}^{\infty} \delta(x) dx}{\lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} \sqrt{\frac{n^2}{\pi}} \exp(-n^2 x^2) dx} \\ &= \lim_{n \rightarrow \infty} 2 \sqrt{\frac{1}{\pi}} \int_0^{\infty} \exp(-y^2) dy \\ &= \lim_{n \rightarrow \infty} 2 \sqrt{\frac{1}{\pi}} \sqrt{\frac{\pi}{4}} \\ &= \lim_{n \rightarrow \infty} 1 \\ &= 1 \end{aligned}$$

2.4 Lorentzian Pulse

$$\begin{aligned} \delta_n(t) &= \frac{n}{\pi} \frac{1}{1+n^2x^2} \\ &= \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} \frac{n}{\pi} \frac{1}{1+n^2x^2} dx \\ &= \lim_{n \rightarrow \infty} \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{1}{1+n^2x^2} dn \\ &= \lim_{u \rightarrow \infty} \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{1}{1+u^2} du \\ &= \lim_{u \rightarrow \infty} \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{1}{1+u^2} du \end{aligned}$$



$$\int \frac{1}{1+u^2} du = \tan^{-1} u$$

(For reference only , but actually you should already know this !!!)

The following illustrate how to calculate this integral

Let $u = \tan \theta$, $du = \sec^2 \theta d\theta$, at $u = \infty$, $\theta = +\frac{\pi}{2}$, at $u = -\infty$, $\theta = -\frac{\pi}{2}$

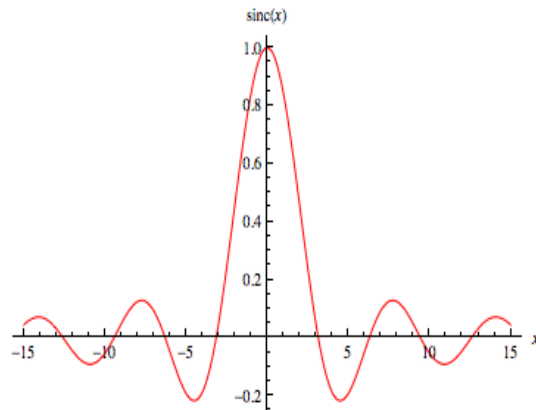
$$\begin{aligned} &\int \frac{1}{1+u^2} du \\ &= \int \frac{1}{1+\tan^2 \theta} \sec^2 \theta d\theta \\ &= \int d\theta \\ &= \theta \\ &= \tan^{-1} u \end{aligned}$$

Thus

$$\begin{aligned} &\lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} \frac{n}{\pi} \frac{1}{1+n^2x^2} dx \\ &= \frac{1}{\pi} \lim_{u \rightarrow \infty} \int_0^{\infty} \frac{1}{1+u^2} du \\ &= \frac{1}{\pi} \lim_{u \rightarrow \infty} [\tan^{-1} u]_{-\infty}^{\infty} \\ &= \frac{1}{\pi} \left[\frac{\pi}{2} - \left(-\frac{\pi}{2} \right) \right] \\ &= 1 \end{aligned}$$

2.5 Sinc Pulse

$$\begin{aligned} \delta_n(t) &= \frac{\sin nx}{\pi x} \\ &= \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} \frac{\sin nx}{\pi x} dx \\ &= \lim_{n \rightarrow \infty} \frac{n}{\pi} \int_{-\infty}^{\infty} \frac{\sin nx}{nx} dx \\ &= \lim_{n \rightarrow \infty} \frac{1}{\pi} \int_{-\infty}^{\infty} \text{sinc} nx dx \\ &= \lim_{u \rightarrow \infty} \frac{1}{\pi} \int_{-\infty}^{\infty} \text{sinc} u du \end{aligned}$$



The integral $\int_{-\infty}^{\infty} \text{sinc} u du$, the Bilateral Sinc integral , BiSi(∞) , has the value of

$$\int_{-\infty}^{\infty} \text{sinc} u du = \pi$$

(All for reference , but actually you should already know this !!!!!)

The following illustrate how to do this integral using Jordan Contour Integral Technique in Complex Analysis

Consider

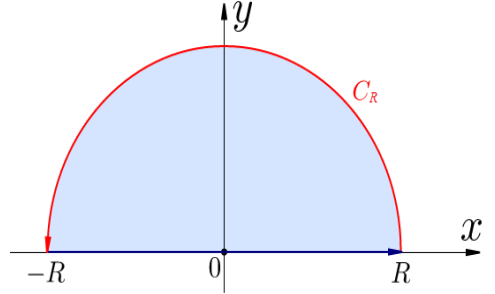
$$\frac{e^{jz}}{z} \text{ has a simple pole at } z = 0$$

To use the Jordan method, construct a function with removable singularity at $z = 0$

$$f(z) = \frac{e^{jz} - 1}{z}$$

Consider Cauchy's integral theorem

$$\oint f(z)dz = \underbrace{\int_{-R}^R f(x)dx}_{\text{along x-axis}} + \underbrace{\int_C f(z)dz}_{\text{circular contour}} = 0$$



So

$$\begin{aligned} & \oint f(z)dz \\ &= \underbrace{\int_{-R}^R f(x)dx}_{\text{along x-axis}} + \underbrace{\int_C f(z)dz}_{\text{circular contour}} \\ &= \int_{-R}^R f(x)dx + \int_C \frac{e^{jz} - 1}{z} dz \\ &= \int_{-R}^R f(x)dx + \int_C \frac{-1}{z} dz + \int_C \frac{e^{jz}}{z} dz \end{aligned}$$

Now consider the $\int_C \frac{1}{z} dz - \int_C \frac{e^{jz}}{z} dz$

For the first integral

$$\begin{aligned} & \int_C \frac{1}{z} dz \\ &= \int_0^\pi e^{-j\theta} j e^{j\theta} d\theta \quad (z = e^{j\theta}) \\ &= \int_0^\pi j d\theta \\ &= j\pi \end{aligned}$$

For the second integral, consider the ML-inequality

$$\left| \int_C \frac{e^{jz}}{z} dz \right| \leq \frac{1}{R} \int_C |e^{jz}| |dz| < \frac{1}{R} \int_0^\pi dz = \frac{\pi}{R}$$

Therefore

$$\begin{aligned} \oint f(z)dz &= \int_{-R}^R f(x)dx - j\pi + \int_C \frac{e^{jz}}{z} dz = 0 \\ \iff \int_{-R}^R f(x)dx - j\pi + \int_C \frac{e^{jz}}{z} dz &= 0 \\ \iff \int_{-R}^R f(x)dx - j\pi &= - \int_C \frac{e^{jz}}{z} dz \\ \iff \left| \int_{-R}^R f(x)dx - j\pi \right| &= \left| \int_C \frac{e^{jz}}{z} dz \right| < \frac{\pi}{R} \\ \iff \left| \int_{-R}^R f(x)dx - j\pi \right| &< \frac{\pi}{R} \end{aligned}$$

(All for reference , but actually you should already know this !!!!!)

Now consider

$$\begin{aligned}
 \operatorname{Im} \left[\int_{-R}^R f(z) dz \right] &= \operatorname{Im} \int_{-R}^R \frac{e^{jz} - 1}{z} dz \\
 &= \int_{-R}^R \frac{\cos z + j \sin z - 1}{z} dz \\
 &= \int_{-R}^R \frac{\sin z}{z} dz \\
 &= \int_{-R}^R f(z) dz
 \end{aligned}$$

Therefore

$$\begin{aligned}
 \left| \int_{-R}^R f(x) dx - j\pi \right| &< \frac{\pi}{R} \\
 \Leftrightarrow \left| \operatorname{Im} \int_{-R}^R f(z) dz - j\pi \right| &< \frac{\pi}{R}
 \end{aligned}$$

Thus the imaginary unit can be taken out

$$\Leftrightarrow \left| \int_{-R}^R f(x) dx - \pi \right| < \frac{\pi}{R}$$

Notice that since absolute value is non-negative

$$0 < \left| \int_{-R}^R \frac{\sin x}{x} dx - \pi \right| < \frac{\pi}{R}$$

Now , to show the middle term is zero, we have to take the limit

$$\lim_{R \rightarrow \infty} \frac{\pi}{R} = 0$$

Thus, by Squeeze Theorem

$$\begin{aligned}
 \lim_{R \rightarrow \infty} \int_{-R}^R \frac{\sin x}{x} dx - \pi &= 0 \\
 \Leftrightarrow \lim_{R \rightarrow \infty} \int_{-R}^R \frac{\sin x}{x} dx &= \int_{-\infty}^{\infty} \frac{\sin x}{x} dx = \pi
 \end{aligned}$$

And therefore

$$\begin{aligned}
 &\int_{-\infty}^{\infty} \delta(x) dx \\
 &= \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} \frac{\sin nx}{\pi x} dx \\
 &= \lim_{u \rightarrow \infty} \frac{1}{\pi} \int_{-\infty}^{\infty} \operatorname{sinc} u du \\
 &= \lim_{n \rightarrow \infty} \frac{1}{\pi} \\
 &= 1
 \end{aligned}$$

—END—